

# Formation of carbon-encapsulated metallic nano-particles from metal acetylides by electron beam irradiation

## Ways for controlled patterning of metallic lines and magnetic dot arrays

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**Abstract.** Transition metal acetylides,  $MC_2$  ( $M=Fe, Co$  and  $Ni$ ), exhibit ferromagnetic behavior of which  $T_C$  is characteristic of their size and structure.  $CoC_2$  synthesized in anhydrous condition exhibited cubic structure with disordered  $C_2^{2-}$  orientation. Once being exposed to water (or air), the particles behave ferromagnetically due to the lengthening of the Co–Co distance by the coordination of water molecules to  $Co^{2+}$  cations. Heating of these particles induces segregation of metallic cores with carbon mantles. Electron beam or 193 nm laser beam can produce nanoparticles with metallic cores covered with carbon mantles.

**PACS.** 36.40.-c Atomic and molecular clusters – 36.40.Cg Electronic and magnetic properties of clusters – 36.40.Mr Spectroscopy and geometrical structure of clusters – 75.50.Tt Fine-particle systems; nanocrystalline materials

## 1 Introduction

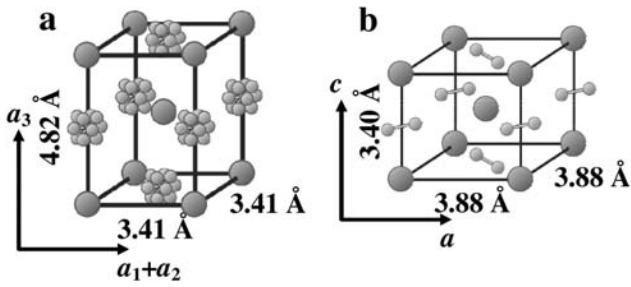
In contrast to the S–S bonds in the transition-metal pyrite compounds ( $MS_2$ ), the  $C\equiv C^{2-}$  bonds of acetylide compounds ( $MC_2$ :  $M=Fe, Co$  and  $Ni$ ) exhibit molecular-nature of the  $C_2^{2-}$  dianion. The orthogonal two  $\pi$ -orbitals of the dianion molecule are expected to strongly interact with  $d$ -orbitals of the transition-metal dication. The infrared absorption spectra of these acetylides show the C=C stretching vibrational bands in doublet in the same region as that of  $CaC_2$ . X-ray absorption near edge spectra of  $CoC_2$  confirmed that the Co atom is in the dication state [1]. The magnetic and electronic properties of the acetylide compounds are very much related to geometrical configuration of metal atoms surrounded by  $C_2^{2-}$  dianions. Close location of the two metal atoms is prone to induce antiferromagnetic (AF) interaction, while ferromagnetic (FM) interaction is also expected for elongated metal-metal distances with the inserted molecule. We reported that  $CoC_2$  is generated from the reaction of  $Co_4(CO)_{12}$  with solvent  $CH_2Cl_2$  not only by heating but also by UV irradiation [1, 2]. In contrast to explosive property of  $M_2C_2$  type acetylides such as  $Ag_2C_2$  and  $Cu_2C_2$ ,  $CoC_2$  and  $NiC_2$  are rather stable in atmospheric condition. Thus, one can make the transition metal acetylides efficiently through the ion exchange reaction of metal chloride with  $CaC_2$  in dehydrated acetonitrile solution. The

$MC_2$  compounds are expected to show  $CaC_2$  type or  $MgC_2$  type structure where metal cations and  $C\equiv C^{2-}$  are alternatively stacked like a rock salt crystal, although the relative orientation of the anion molecule may depend on the synthetic condition. Closer location of the two metal atoms may induce AF interaction, while the elongation of the metal-metal distances may result in the FM interaction due to the elongated rhombus configuration of the  $M-C_2-M$  four-center unit. We demonstrate this property for  $CoC_2$  nanocrystals synthesized at 78 °C. High temperature treatment of  $FeC_2$ ,  $CoC_2$ , and  $NiC_2$ , however, induces charge neutralization resulting in the segregation of metal cores with graphitic or pyrolytic carbon mantles [3]. This phenomenon can be utilized for electron beam manufacturing of nano-dots or nano-wires on a Si substrate. VUV laser beam processing is also possible for acetylide thin layers.

## 2 Water induced structural change and ferromagnetism of $CoC_2$

Owing to the highly absorbency of  $CoC_2$ , ca. 1–2  $H_2O$  molecules are coordinated to the  $Co^{2+}$  dication when  $CoC_2$  is exposed to air or water. To investigate the influence of the coordinated water on the structure and magnetic property of  $CoC_2$ , XRD and SQUID measurement were done for hydrous and anhydrous  $CoC_2$ . Hydrous  $CoC_2$  was prepared by the air exposure of

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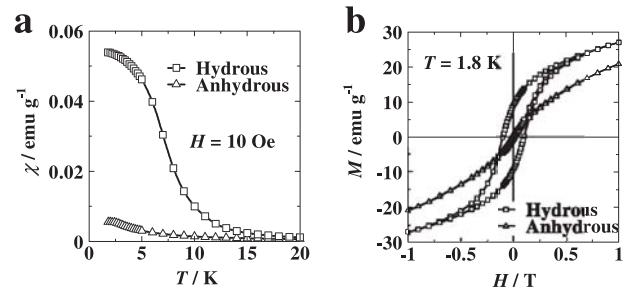


**Fig. 1.** The crystal structure of anhydrous (a) and hydrous (b)  $\text{CoC}_2$ . The figure of the anhydrous  $\text{CoC}_2$  is a part of the fcc-unit cell. Large and small gray circles indicate cobalt and carbon atoms, respectively. The disorder of the  $\text{C}_2^{2-}$  dianions in anhydrous  $\text{CoC}_2$  is depicted as a static disorder on several carbon positions.

anhydrous  $\text{CoC}_2$  for 2 hours. Though the washing with water also gives hydrous  $\text{CoC}_2$  in a short time, air-exposed sample is suitable to measure the magnetic property thanks to its larger structural domain size.

The observed XRD pattern of the as-prepared anhydrous  $\text{CoC}_2$ , which is mostly originated from the scattering at the  $\text{Co}^{2+}$  dications, suggests the fcc-sublattice of  $\text{Co}^{2+}$  with the lattice constants  $a_1 = a_2 = a_3 = 4.82 \text{ \AA}$ . Judging from the crystal structures of  $\text{CaC}_2$  and  $\text{MgC}_2$ , we expect that the  $\text{Co}^{2+}$  and  $\text{C}_2^{2-}$  ions form NaCl-like alternate stack, where the fcc arrangement of the  $\text{Co}^{2+}$  dications indicates the (static) orientation disorder of the  $\text{C}_2^{2-}$  dianions as shown in Figure 1a. The cubic structure involving the orientation disorder of  $\text{C}_2^{2-}$  dianions is also observed in  $\text{CaC}_2$  at high temperature [4]. In addition, the small peak widths ( $\Delta 2\theta \sim 2.00 \times 10^{-3} \text{ rad}$  at  $\theta = 24.0^\circ$ ,  $\lambda = 0.9988 \text{ \AA}$ ) imply the large domain size  $D$  estimated as  $D > 45 \text{ nm}$  by using Scherrer's equation.

Although  $\text{CoC}_2$  is water-stable material unlike  $\text{CaC}_2$  and  $\text{MgC}_2$ , absorbed water changes the crystal structure of  $\text{CoC}_2$  drastically. After air-exposure, the XRD peaks become broad and the positions change as  $\text{Co}^{2+}$  cations form the body-centered-tetragonal sublattice with the lattice constants  $a = b = 3.88$ ,  $c = 3.40 \text{ \AA}$ . These values are consistent with the crystal structure forecasted from the EXAFS results ( $a = b = 3.85 \text{ \AA}$  and  $c = 3.61 \text{ \AA}$ ), where the crystal structure of  $\text{CoC}_2$  is regarded as the  $\text{MgC}_2$ -type structure [5] as shown in Figure 1b. In this phase, the  $a$ - and  $b$ -axes are expanded due to the absorbed waters coordinated to the  $\text{Co}^{2+}$ , while the  $c$ -axis becomes shorter. The orientation of the  $\text{C}_2^{2-}$  dianions are also affected by the lattice expansion and/or steric hindrance of the water, resulting in the structural change from the isotropic disordered orientation (Fig. 1a) to anisotropic ordered orientation (Fig. 1b). Because the structural change is too drastic to be done in concert, a  $\text{CoC}_2$  particle is divided into many small domains. Indeed, the XRD peak widths of the hydrous  $\text{CoC}_2$   $\Delta\theta \sim 1.25 \times 10^{-2} \text{ rad}$  at  $2\theta = 34.6^\circ$ ,  $\lambda = 1.5418 \text{ \AA}$  is remarkably larger than that of anhydride, suggesting the small domain size  $D \sim 10 \text{ nm}$  of the hydrous  $\text{CoC}_2$ .

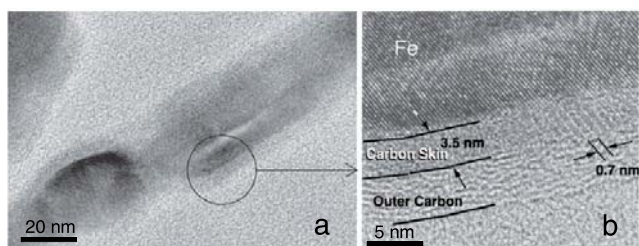


**Fig. 2.** Field-cooled magnetic susceptibility (a) and magnetization curves (b) of the anhydrous and hydrous  $\text{CoC}_2$ .

Inter-molecular magnetic interaction is sensitive to the arrangement of the molecules. Therefore, it is expected that the magnetism of the hydrous and anhydrous  $\text{CoC}_2$  are largely different. Figure 2a shows the temperature dependence of the field-cooled magnetic susceptibilities  $\chi$  of  $\text{CoC}_2$  before and after the air-exposure measured at  $H = 10 \text{ Oe}$ . It is noted that ICP measurement reveals the  $\text{CoC}_2$  containing ca. 30 at% of starting material ( $\text{CaC}_2$ ), therefore observed  $\chi$  is probably ca. 30% smaller than that of pure  $\text{CoC}_2$ .  $\chi$  of the anhydrous  $\text{CoC}_2$  obeys the Curie-Weiss law in the high temperature range above ca. 70 K with Curie constant  $C = 1.1 \text{ emu K/mol}$  and AF Weiss temperature  $\Theta = -10 \text{ K}$ , the former of which is obviously larger than the value of  $\text{Co}^{2+}$  cation ( $C = 0.375$ ), suggesting the short-range strong FM interaction. Below 2.5 K, the anhydrous  $\text{CoC}_2$  shows ferromagnetism as discussed in later, but the FM part is only less than 8%.

Although  $\chi$  of the hydrous  $\text{CoC}_2$  also obeys the Curie-Weiss law above ca. 100 K,  $C = 1.5 \text{ emu K/mol}$  is significantly larger than that of anhydrous  $\text{CoC}_2$  suggesting the expansion of the FM domain, while AF  $\Theta = -3 \text{ K}$  indicates the weakening of the AF interaction between FM domains. The difference of the magnetism between hydrous and anhydrous  $\text{CoC}_2$  is evidently caused by the water-induced structural changes; that is, the orientation ordering of  $\text{C}_2^{2-}$  dianions and the change of the inter- and intra-chain distances. Judging from the fact that the interaction is weakened by lengthening the inter-molecular distance, we conclude that the AF and FM interactions are attributed to the inter- and intra-chain interaction. In the anhydrous phase, the orientation of  $\text{C}_2^{2-}$  dianions are disordered, resulting in the fragmental small FM domains owing to the disordered FM and AF interactions. On the other hand, the orientation of the  $\text{C}_2^{2-}$  dianions tends to be ordered in hydrous phase, where the FM interactions are connected to each other and form the FM chain elongated parallel to the  $c$ -axis, resulting in the large FM domains.

The water-induced ferromagnetism is clearly observed in the magnetization curves as shown in Figure 2b. In the case of anhydrous  $\text{CoC}_2$ , only a small amount of the FM component shows hysteresis loop with coercive force  $H_C = 150 \text{ Oe}$  and remanent magnetization  $M_{\text{REM}} = 0.7 \text{ emu g}^{-1}$ , while the greater part of  $\text{CoC}_2$  behaves as paramagnet with weak AF interaction. The small amount of the FM component is perhaps caused by the residual water in the resin used for the SQUID measurement.



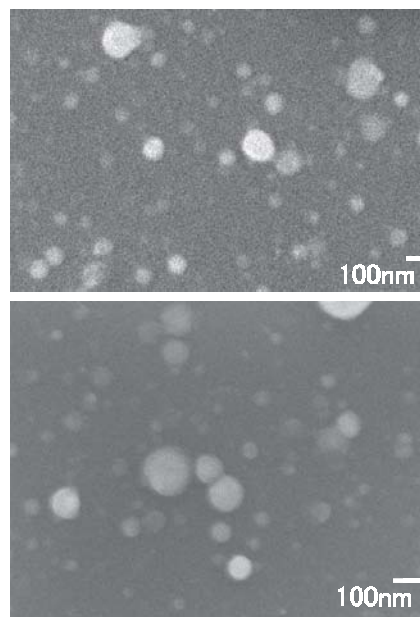
**Fig. 3.** TEM image of the Fe nano-rod (a) and its enlargement (b). Iron core is covered with the graphite skin whose thickness is ca. 3.5 nm (see text).

Contrary to the anhydrous  $\text{CoC}_2$ , the hysteresis loop of the hydrous  $\text{CoC}_2$  is clearly observed with  $H_c = 1000$  Oe and  $M_{\text{REM}} = 9.0$  emu  $\text{g}^{-1}$ . The results clearly show the importance of the water in the ferromagnetism of  $\text{CoC}_2$ .

Here, we briefly note the effect of the water on  $\text{FeC}_2$  and  $\text{NiC}_2$ . Anhydrous  $\text{FeC}_2$  shows paramagnetism except a very small ferromagnetic part, as like as anhydrous  $\text{CoC}_2$ . However, different from  $\text{CoC}_2$ , water-absorption decomposes  $\text{FeC}_2$  into acetylene and iron hydroxide like as  $\text{CaC}_2$ . In the case of  $\text{NiC}_2$ , the major part shows ferromagnetism even in anhydrous phase. Furthermore, water-absorption does not occur (or very slow) in  $\text{NiC}_2$ . The difference may be originated from the configuration of the  $d$ -electrons, but detail is still unclear.

### 3 Carbon-skinned Fe rods

By the rapid stirring of the acetonitrile solution with  $\text{FeCl}_2$  and  $\text{CaC}_2$  fine particles, one can get the Fe nanocrystals with graphitic carbon-skins. Carbon-protected metal nanoparticle is a promising material for magnetic devices owing to its chemical stability [6]. The reaction time should be as long as 48 hours at 250 °C. The intermediate product of  $\text{FeC}_2$  shows charge neutralization reaction and segregation into metal and carbon components. Though bulk carbon and bulk iron are more stable than  $\text{FeC}_2$ , alternate stack of  $\text{Fe}^{2+}$  and  $\text{C}_2^{2-}$  forbids the neutralization at low temperature, where the neutral  $\text{Fe-C}_2$  pair caused by electron transfer is re-ionized by Coulomb potential of adjacent ions. On the contrary to low temperature, neutral iron and carbon atoms can easily migrate and gather before re-ionization at high temperature. Once the “iron core” is formed, the core collects neighbouring iron atoms and becomes iron nanoparticle covered with residual carbon atoms. X-ray diffraction pattern of the products exhibits an asymmetric broad band at  $2\theta = 26.0^\circ$  that was assigned to graphitic carbon. From the Scherrer equation, the average size of the iron core is estimated as 30 nm. Figure 3 shows the TEM images of a nano-rod with a neck. Figure 3b is an expanded image of the outer area. Although most of nano-crystals with short lengths exhibit single carbon-skin with a thickness of 3.5 nm, this long rod wears double carbon layers. The first skin layer is directly bonded with an iron metal core and the rod wears the second outer-mantle with planes parallel to the iron surface. The direct bonding of the graphitic axis to Fe axes



**Fig. 4.** Typical SEM images of the carbon-encapsulated Co nanoparticles formed by 193 nm laser irradiation on  $\text{CoC}_2$  in Ar atmosphere (0.5 torr).

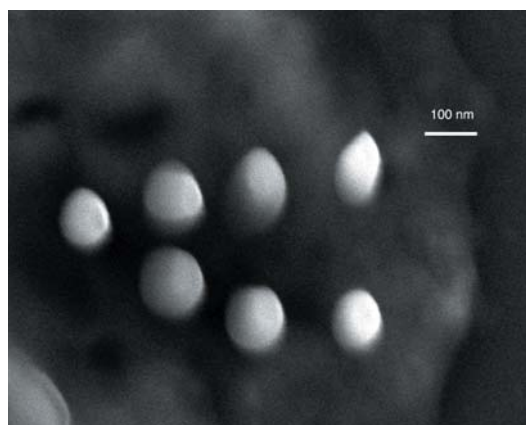
was attributed to accidental coincidence of the Fe–Fe distance (2.866 Å) with the  $\text{C}_1$ – $\text{C}_4$  distance of a graphite hexagonal ring (2.842 Å). Due to the modification of the surface structure, the saturation magnetization at 300 K of the carbon-skinned Fe nano-particles with an average diameter of 60 nm was as high as 80% of that at 1.8 K [3].

The carbon-skinned metal nanoparticles of Co and Ni can also be obtained from acetylide, but the “skin” is amorphous carbon due to the lattice mismatch between graphite and Ni or Co metals.

The second ionization potentials of Ni, Co, and Fe are 18.1688, 17.083, and 16.1878 eV, respectively, and the first ionization potentials are 7.6398, 7.8810, and 7.9024 eV, respectively [7].

### 4 Carbon-encapsulated Co by VUV irradiation of $\text{CoC}_2$

Because of its water- and air-stability,  $\text{CoC}_2$  is suitable to create a thin film of carbon-encapsulated metal nanoparticles by laser irradiation, where we must treat a bulk amount of the material.  $\text{CoC}_2$  disk is irradiated by a 193 nm ArF laser beam in Ar atmosphere. The particles sputtered from the surface were deposited on a Si(111) surface. Figure 4 shows typical SEM images of the surface. White particles are somewhat charged up due to the insulation by the carbon mantle of the spheric particles with Co cores. EDS analysis reveals that the particles consist of carbon and cobalt atoms. These figures demonstrate that 193 nm laser beam irradiation produces carbon-encapsulated spherical Co particles. The size of the particles are dependent on the laser power. Higher intensity causes larger particles. For the generation of the



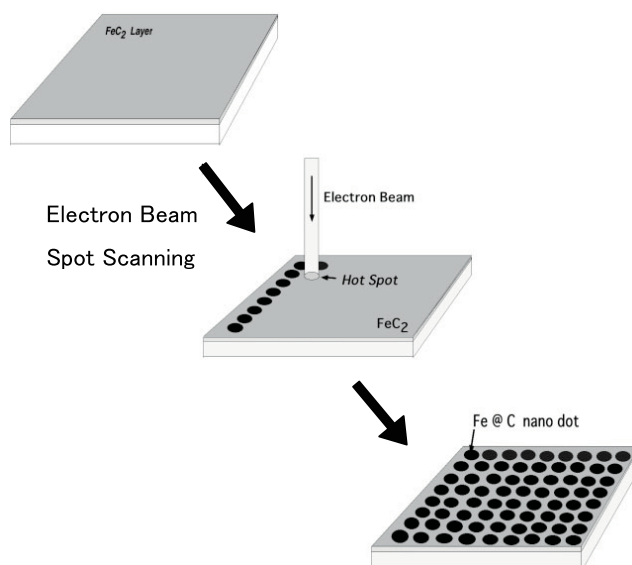
**Fig. 5.** Nickel nanodots from  $\text{NiC}_2$  formed by the electron beam irradiation.

particles in Figure 4 we controlled the laser power as high as  $200 \text{ mJ/cm}^2$ .

## 5 Nano-dots formation by electron beam heating of $\text{NiC}_2$

Spot irradiation of weak electron beams can generate the metallic dots covered with carbons. Figure 5 shows an example of the preparation of nickel particles with carbon mantles with a scanning electron microscope (SEM), where we used  $\text{NiC}_2$  because of its lower segregation temperature ( $150 \text{ }^\circ\text{C}$ ) than that of  $\text{CoC}_2$  or  $\text{FeC}_2$  ( $> 200 \text{ }^\circ\text{C}$ ). The lower segregation temperature is owing to the large second ionization potentials of Ni [7], and is suited to create a dot by weak electron beam. In the creation of dot,  $\text{NiC}_2$  flat particles were situated on a  $\text{Si}(100)$  surface. Electron beam current was adjusted to  $8 \text{ pA}$  and irradiation was performed at arbitrary points. The size of the white spots became larger with elongating irradiation time and grew up to  $100 \text{ nm}$  after 3 minutes. The shape of the spots reflects the flatness of the surface. Appearance of white spots suggests that the carbon mantles on the metallic cores are very poor electric conductors and charges are localized on the surface of the amorphous carbon mantle. The observation of the energy dispersive X-ray spectrum suggested that the yield of carbon signals relative to that of Ni increases with increasing irradiation time. This indicates that the metallic core as well as the thickness of the carbon mantle grows with elongated heating time.  $\text{NiC}_2$ , in particular, is easily segregated even during the SEM image observations when electron beam current is not weakened.

This phenomenon can be applied for making nano-dots arrays embedded in an acetylide matrix. Figure 6 illustrates how to make nano-dots arrays on a silicon surface. Acetylide thin layer can be made in diluted Ar atmosphere by  $193 \text{ nm}$  laser sputtering at a moderate photon flux. The



**Fig. 6.** Schematic illustration of the nano-dots arrays formation by electron beam heating.

acetylide layer is composed of minute particles with diameters of sub-nanometer or nanometers. These tiny acetylide particles are superparamagnetic and, therefore, they can shield the magnetic moment of the magnetic dots. This role is very important, since the magnet-magnet interaction between the big dots is so large that the memory written in each dot can be easily destroyed through this interaction. The insulation by the surrounding tiny paramagnets from the strong inter-magnet interaction is quite necessary for the stabilization of the magnetic moments of the dots. For the dots larger than  $200 \text{ nm}$ , the direct heating can be done with focused  $193 \text{ nm}$  beam through a mask pattern. Carbon-covered Co spherical-dots with diameters of order of several hundred nm were formed with focused  $193 \text{ nm}$  laser irradiation on  $\text{CoC}_2$  disks. This indicates that laser irradiation can melt  $\text{CoC}_2$  and segregate Co cores as aligned magnets.

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