

# **BOREL-TYPE SUMMABILITY METHODS**

**THESIS SUBMITTED FOR THE DEGREE OF  
DOCTOR OF PHILOSOPHY**

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## Chapter-I

### INTRODUCTION

Summability transformations help us to generalize the concept of limit of a sequence or series, and thus provide us a method to assign limits even to sequences which are divergent. These transformations or methods can be classified into two types.

- (i) Sequence to sequence transformations
- (ii) Sequence to function transformations

Sequence to sequence transformations are accomplished using infinite matrices. Consider an infinite matrix  $C = (c_{nk})$  and a sequence  $\{s_n\}$ ,  $n = 0, 1, 2, \dots$ . Form the new sequence  $\{t_n\}$  defined by

$$t_n = \sum_{k=0}^{\infty} c_{nk} s_k$$

We shall assume that the series converges for every  $n$ .

$\{t_n\}$  is called the  $C$ -transform of the given sequence  $\{s_n\}$ . If  $\{t_n\}$  converges to  $t$ , then  $t$  is called the  $C$ -limit of  $\{s_n\}$  and we write  $s_n \rightarrow t(C)$ .

A transformation  $C$  is called a regular summability transformation if it preserves limit in the case of convergent sequences. That is

$$s_n \longrightarrow s \implies t_n \longrightarrow s(C)$$

Silverman-Toeplitz theorem gives necessary and sufficient conditions for a matrix to represent a regular method and thus help us to construct regular transformations. This theorem can be stated as: The necessary and sufficient conditions that the matrix  $C = (c_{nk})$  represents a regular transformation are:

- (i)  $\sum_{k=0}^{\infty} |c_{nk}| < M$ , for some  $M$  and for all  $n=0,1,2,\dots$  .
- (ii)  $\lim_{n \rightarrow \infty} c_{nk} = 0$ , for each  $k = 0,1,2,\dots$  .
- (iii)  $\lim_{n \rightarrow \infty} \sum_{k=0}^{\infty} c_{nk} = 1$

As an example for the sequence to sequence transformation consider the  $(C,1)$  mean (Cesaro mean of order 1). The transformed sequence  $\{t_n\}$  of a given sequence  $\{s_n\}$  is defined by

$$t_n = \frac{s_0 + s_1 + \dots + s_{n-1}}{n}$$

If  $\{s_n\} = 1, 0, 1, 0, \dots$ , then  $\{s_n\}$  is not convergent where as  $s_n \longrightarrow \frac{1}{2}(C, 1)$ . The matrix of  $(C, 1)$  is given by

$$\begin{bmatrix} 1 & 0 & 0 & 0 & \dots \\ \frac{1}{2} & \frac{1}{2} & 0 & 0 & \dots \\ \frac{1}{3} & \frac{1}{3} & \frac{1}{3} & 0 & \dots \\ \dots & \dots & \dots & \dots & \dots \\ \dots & \dots & \dots & \dots & \dots \end{bmatrix}$$

As an example for sequence to function transformation, consider the Abel method, defined by

$$s_n \longrightarrow s(A),$$

if,

$$\lim_{r \rightarrow 1-} (1-r) \sum_{n=0}^{\infty} s_n r^n = s$$

If  $\{s_n\}$  denotes the sequence of partial sums of the series  $\sum_0^{\infty} a_n$ , then we have the relation

$$\frac{\sum a_n r^n}{1-r} = \sum_0^{\infty} s_n r^n$$

Hence for a series  $\sum_0^{\infty} a_n$ , Abel method is defined as

$$\sum a_n = s(A),$$

if

$$\lim_{r \rightarrow 1-} \sum_{n=0}^{\infty} a_n r^n = s$$

In particular the series  $1-2+3-4+ \dots$  is summable Abel with limit  $\frac{1}{4}$ . The partial sums of the above series determine the sequence  $1, -1, 2, -2, \dots$  which is not (C,1) summable.

In the field of summability theory, Tauberian theorems occupy an important position. These theorems provide results of the following type. If a sequence  $\{s_n\}$  is summable by a method C and also  $s_n$  satisfy some condition (called Tauberian condition) then  $\{s_n\}$  is convergent. The first theorem of this character was proved by A. Tauber in 1897 and he proved the following result ([16], Theorem 85).

" If  $\sum a_n$  is summable (Abel) to  $s$  and  $a_n = o(\frac{1}{n})$ , then  $\sum a_n$  converges to  $s$ ". This theorem has been generalized by showing that the result is true even when  $a_n = O(\frac{1}{n})$  ([16], Theorem 90). A similar theorem for Borel-summability can be stated as follows:

" If  $\sum a_n$  is Borel summable to  $s$  and  $a_n = O(\frac{1}{\sqrt{n}})$ , then  $\sum a_n$  converges to  $s$ ". ([16], Theorem, 156).

Definitions, properties and theorems concerning a large number of regular transformations can be obtained from the book "Divergent Series" by G.H. Hardy ([16]).

Another variant of Tauberian theorem prove results of the following character. " If  $\rho \geq -\frac{1}{2}$ ,  $a_n = o(n^\rho)$  and  $\Sigma a_n$  is summable Borel to  $s$ , then  $\Sigma a_n$  is summable  $(C, 2\rho+1)$  to  $s$ ". ([16], Theorem 147). Borel method is a particular case of Borel-type method  $(B, \alpha, \beta)$ . (Results on  $(B, \alpha, \beta)$  method form the material of this thesis). The above theorem was generalized to  $(B, \alpha, \beta)$  method by Borwein, D [5]. This was again improved by Kwee, B [25] by showing that the result is true with  $a_n = O(n^\rho)$ . A few papers dealing results of this nature are [7], [8], [10], [26].

The study of Gibbs phenomenon and Lebesgue constants for different summability transformations had been undertaken by many researchers ([18],[19],[27],[28],[29],[30]). Summability transformations help us to extend the domain of convergence of a series of functions. The domain of convergence of a series of Legendre polynomials for different summability transformations had been investigated by

many authors ([11], [21], [22], [32]). Tauberian constants for many summability methods had been determined. The problem generally considered in this area can be stated in the following way. Let  $\{s_n\}$  be the sequence of partial sums of  $\sum a_n$  and assume  $a_n$  satisfies some tauberian condition. Let  $T(x) = \sum_{n=0}^{\infty} c_n(x)s_n$  be a sequence to function summability transformation. Then the results give estimates of

$$\lim_{n \rightarrow \infty} \sup_{x_n \rightarrow \infty} |T(x_n) - s_n|$$

when neither  $\lim T(x)$  nor  $\lim s_n$  is assumed to exist. [1], [20], [31], [33] deal results of this nature.

A brief introduction to Summability transformations touching all aspects of the theory is almost an impossible task. Hence in the above introduction, I have restricted the concepts to those which are relevant to the topics discussed in this thesis.

### Brief summary of the results in the thesis

Before summarising the results, we first define the Borel-type transformation  $(B, \alpha, \beta)$  which is a generalization of the classical Borel transform. After Borwein, D([5]) we may define  $(B, \alpha, \beta)$  summability as follows:



Let  $\{s_n\}$ ,  $n=0,1,2,\dots$  be a sequence of real or complex numbers. Suppose that  $\alpha > 0$ ,  $\beta$  is real and  $N$ , a non-negative integer such that  $\alpha N + \beta > 0$ . The sequence  $\{s_n\}$  is said to be  $(B, \alpha, \beta)$  summable to  $s$ , if

$$\lim_{x \rightarrow \infty} \alpha e^{-x} \sum_{n=N}^{\infty} \frac{x^{\alpha n + \beta - 1}}{\Gamma(\alpha n + \beta)} s_n = s$$

The  $(B, \alpha, \beta)$  method is regular and reduces to the classical Borel method when  $\alpha = \beta = 1$ . Being a generalization of the Borel transform  $(B, \alpha, \beta)$  method is also called Borel-type summability method.

In chapter 2, the study of Gibbs phenomenon with regard to  $(B, \alpha, \beta)$  summability is undertaken. It is shown that the Borel-type summability completely preserves the Gibbs phenomenon for Fourier series.

In chapter 3, the Lebesgue constants for  $(B, \alpha, \beta)$  method is calculated. It is shown that the Lebesgue constants  $L_B(x)$  for  $(B, \alpha, \beta)$  method is given by

$$L_B(x) = \frac{2}{\pi} \log \left( \frac{2x}{\pi} \right) - \frac{2}{\pi} C - \frac{2}{\pi} \int_0^{\pi} \psi \left( \frac{t}{\pi} \right) \sin t \, dt + O \left( \frac{1}{\sqrt{x}} \right) \quad (x \rightarrow \infty)$$

where  $C$  is the Euler-Mascheroni constant and

$$\Psi(t/\pi) = \frac{\Gamma'(t/\pi)}{\Gamma(t/\pi)}$$

In chapter 4, the domain of summability of Legendre polynomials by the  $(B, \alpha, \beta)$  transform is obtained. It is shown that the sequence  $\{s_k(z, w)\}$  of partial sums of the series of Legendre polynomials  $\sum_{n=0}^{\infty} (2n+1) P_n(z) Q_n(w)$  is summable  $(B, \alpha, \beta)$  to  $(w-z)^{-1}$  in the region

$$\left\{ z : \operatorname{Re} \left\{ \frac{\xi(\phi)}{\tau(\psi)} \right\}^{1/\alpha} < \lambda, \left| \frac{\xi(\phi)}{\tau(\psi)} \right| < M, 0 \leq \phi \leq \pi, 0 \leq \psi \right\}$$

where  $M$  is a positive constant and  $0 < \lambda < 1$ .

In chapter 5, we prove a result on Tauberian constants for  $(B, \alpha, \beta)$  transform of a series  $\sum u_k$  with the condition  $\limsup | \sqrt{k} u_k | = L < \infty$ . It is shown that if  $m \rightarrow \infty$ ,  $t \rightarrow \infty$  such that

$$\limsup \frac{m-t}{\sqrt{t}} = Q < \infty, \text{ then}$$

$$\limsup |B(t) - s_m| < A.L, \text{ where}$$

$$A = \sqrt{\frac{2\alpha}{\pi}} Q \int_0^Q e^{-\frac{\alpha}{2} z^2} dz + \sqrt{\frac{2}{\alpha\pi}} e^{-\frac{\alpha}{2} Q^2} Q^2$$

## Chapter-II

### GIBBS PHENOMENON FOR $(B, \alpha, \beta)$ SUMMABILITY

#### Preliminaries

In this section we define the Borel-type summability transformations. We also list some of the basic inequalities and estimates satisfied by the Borel-type transform which will be used in the sequel. The proof of these inequalities and estimates can be found in [5]. These are generalization of the corresponding results for the Borel-transform (cf [16], theorem, 137).

#### Definition of the Borel-type transforms $(B, \alpha, \beta)$

Let  $\{s_n\}$ ,  $n=0,1,2,\dots$  be a sequence of real or complex numbers. Suppose that  $\alpha > 0$ ,  $\beta$  is real and  $N$  a non-negative integer such that  $\alpha N + \beta > 0$ . The sequence  $\{s_n\}$  is said to be  $(B, \alpha, \beta)$  summable to the sum  $s$ , if

$$\text{Limit}_{x \rightarrow \infty} \alpha e^{-x} \sum_{n=N}^{\infty} \frac{x^{\alpha n + \beta - 1}}{\Gamma(\alpha n + \beta)} s_n = s.$$

The  $(B, \alpha, \beta)$  method is regular and reduces to the classical Borel method when  $\alpha = \beta = 1$ .

Lemma 2.1. (cf [5])

$$\alpha e^{-x} \sum_{n=N}^{\infty} \frac{x^{\alpha n + \beta - 1}}{\Gamma(\alpha n + \beta)} = 1 + o(1) \quad (x \rightarrow \infty).$$

Lemma 2.2. (cf [5])

$$\text{Let } x > 0, \quad 0 < \delta < \frac{1}{\alpha}, \quad \frac{1}{2} < \xi < \frac{2}{3}, \quad \gamma = \frac{1}{3}(\alpha\delta)^2$$

$0 < \eta < 2\xi - 1$  and let

$$u_n = u_n(x) = \alpha e^{-x} \frac{x^{\alpha n + \beta - 1}}{\Gamma(\alpha n + \beta)}, \quad n = N, N+1, \dots$$

Then,

- (i)  $\sum_{n=N}^{\infty} u_n \rightarrow 1$  as  $x \rightarrow \infty$
- (ii)  $u_n \leq u_{n+1}$  when  $n \leq \frac{x}{\alpha} - \frac{\beta}{\alpha} - 1$ , and  
 $u_{n+1} \leq u_n$  when  $n \geq \frac{x}{\alpha} - \frac{\beta}{\alpha} + \frac{1}{\alpha}$
- (iii)  $\sum_{|n - \frac{x}{\alpha}| > \delta x} u_n = O(e^{-\gamma x}) \quad (x \rightarrow \infty)$
- (iv)  $\sum_{|n - \frac{x}{\alpha}| > x^\xi} u_n = O(e^{-x^\eta}) \quad (x \rightarrow \infty)$

$$\begin{aligned}
 \text{(v)} \quad u_n &= \frac{\alpha}{\sqrt{2\pi x}} e^{-\frac{\alpha^2 h^2}{2x}} \left[ 1 + o\left(\frac{|h|+1}{x}\right) + o\left(\frac{|h|^3}{x^2}\right) \right] \\
 &= \frac{\alpha}{\sqrt{2\pi x}} e^{-\frac{\alpha^2 h^2}{2x}} \left[ 1 + o(x^{3\xi-2}) \right] \text{ where } h = n - \frac{x}{\alpha}
 \end{aligned}$$

(vi) If  $\theta > 0$  fixed, then

$$\sum_{|n - \frac{x}{\alpha}| > \theta x^\xi} u_n = o(e^{-x^\eta}) \quad (x \rightarrow \infty)$$

$$\text{(vii)} \quad \sum_{|n - \frac{x}{\alpha}| > \lambda \sqrt{x}} u_n < \epsilon, \text{ for } x > x_0(\epsilon), \lambda > \lambda_0(\epsilon)$$

Theorem 2.1. (cf [13])

The  $(B, \alpha, \beta)$  Summability completely preserves the Gibbs phenomenon for Fourier series.

Proof:

Consider the function  $f(t)$  defined on  $[0, 2\pi)$  by

$$\begin{aligned}
 f(t) &= \frac{1}{2}(\pi - t), \quad (0 < t < 2\pi) \\
 &= 0 \text{ if } t = 0
 \end{aligned}$$

and extended outside by periodicity.

The Fourier series of  $f$  is given by

$$\sum_{n=1}^{\infty} \frac{\sin nt}{n}$$

Clearly  $f$  is odd and has a jump  $\pi$  at 0. The sequence  $\{s_n(t)\}$  of partial sums of the above Fourier series is given by

$$s_n(t) = -\frac{t}{2} + \int_0^t \frac{\sin(n + \frac{1}{2})u}{2 \sin \frac{u}{2}} du$$

Let  $B_x(t)$  denote the  $(B, \alpha, \beta)$  transform of  $\{s_n(t)\}$ .

Then,

$$\begin{aligned} B_x(t) &= \alpha e^{-x} \sum_{n=N}^{\infty} \frac{x^{\alpha n + \beta - 1}}{\Gamma(\alpha n + \beta)} s_n(t) \\ &= \alpha e^{-x} \sum_{n=N}^{\infty} \frac{x^{\alpha n + \beta - 1}}{\Gamma(\alpha n + \beta)} \left( -\frac{t}{2} + \int_0^t \frac{\sin(n + \frac{1}{2})u}{2 \sin \frac{u}{2}} du \right) \\ &= -\frac{t}{2} \alpha e^{-x} \sum_{n=N}^{\infty} \frac{x^{\alpha n + \beta - 1}}{\Gamma(\alpha n + \beta)} + \alpha e^{-x} \sum_{n=N}^{\infty} \left( \int_0^t \frac{\sin(n + \frac{1}{2})u}{2 \sin \frac{u}{2}} du \right) \frac{x^{\alpha n + \beta - 1}}{\Gamma(\alpha n + \beta)} \\ &= \Sigma_1 + \Sigma_2 \end{aligned} \tag{2.1}$$

say. Applying Lemma 2.1, it follows that

$$\Sigma_1 = \frac{-t}{2} [1+o(1)] \quad (x \rightarrow \infty) \quad (2.2)$$

By changing the order of integration and summation in  $\Sigma_2$ , which is permissible because of uniform convergence, we obtain.

$$\begin{aligned} \Sigma_2 &= \alpha e^{-x} \int_0^t \left[ \sum_{n=N}^{\infty} \frac{(\sin n + \frac{1}{2})u}{2 \sin \frac{u}{2}} \cdot \frac{x^{\alpha n + \beta - 1}}{\Gamma(\alpha n + \beta)} \right] du \\ &= \alpha e^{-x} \int_0^t \operatorname{Im} \left[ e^{i \frac{u}{2}} \sum_{n=N}^{\infty} e^{inu} \frac{x^{\alpha n + \beta - 1}}{\Gamma(\alpha n + \beta)} \right] \frac{du}{2 \sin \frac{u}{2}} \\ &\quad (\operatorname{Im} \text{ means imaginary part}) \\ &= \alpha e^{-x} \int_0^t \operatorname{Im} \left[ e^{iu(\frac{1}{2} + \frac{1-\beta}{\alpha})} \sum_{n=N}^{\infty} \frac{(xe^{\frac{i u}{\alpha}})^{\alpha n + \beta - 1}}{\Gamma(\alpha n + \beta)} \right] \frac{du}{2 \sin \frac{u}{2}} \end{aligned}$$

Now by using Lemma 2.1, we have

$$\begin{aligned} \Sigma_2 &= \alpha e^{-x} \int_0^t \operatorname{Im} \left[ e^{iu(\frac{1}{2} + \frac{1-\beta}{\alpha})} \frac{1}{\alpha} e^{xe^{\frac{i u}{\alpha}}} \{1+o(1)\} \right] \frac{du}{2 \sin \frac{u}{2}} \\ &= \int_0^t \operatorname{Im} \left[ e^{iu(\frac{1}{2} + \frac{1-\beta}{\alpha})} e^{-x(1-e^{\frac{i u}{\alpha}})} \{1+o(1)\} \right] \frac{du}{2 \sin \frac{u}{2}} \quad (2.3) \end{aligned}$$

Now,

$$\operatorname{Im} \left[ e^{iu(\frac{1}{2} + \frac{1-\beta}{\alpha})} e^{-x(1-e^{\frac{i u}{\alpha}})} \right]$$

$$\begin{aligned}
&= \operatorname{Im} \left[ e^{iu\left(\frac{1}{2} + \frac{1-\beta}{\alpha}\right) - x(1 - \cos \frac{u}{\alpha} - i \sin \frac{u}{\alpha})} \right] \\
&= e^{-x(1 - \cos \frac{u}{\alpha})} \operatorname{Im} \left[ e^{i\left(x \sin \frac{u}{\alpha} + \left(\frac{1}{2} + \frac{1-\beta}{\alpha}\right)u\right)} \right] \\
&= e^{-2x \sin^2 \frac{u}{2\alpha}} \sin\left[x \sin \frac{u}{\alpha} + \left(\frac{1}{2} + \frac{1-\beta}{\alpha}\right)u\right]
\end{aligned}$$

Thus from (2.3) we obtain

$$\begin{aligned}
\Sigma_2 &= \int_0^t e^{-2x \sin^2 \frac{u}{2\alpha}} \sin\left[x \sin \frac{u}{\alpha} + u\left(\frac{1}{2} + \frac{1-\beta}{\alpha}\right)\right] (1+o(1)) \frac{du}{2 \sin \frac{u}{2}} \\
&= \int_0^t e^{-2x \sin^2 \frac{u}{2\alpha}} \sin\left[x \sin \frac{u}{\alpha} + \left(\frac{1}{2} + \frac{1-\beta}{\alpha}\right)u\right] \frac{du}{2 \sin \frac{u}{2}} \\
&= \frac{t}{2\alpha} \int_0^{\frac{t}{2\alpha}} e^{-2x \sin^2 w} \sin\left[x \sin 2w + w(\alpha+2-2\beta)\right] \frac{\alpha dw}{\sin(\alpha w)} \quad (2.4)
\end{aligned}$$

Consequently for small values  $t_x$  of  $t$ , we obtain from (2.1), (2.2) and (2.4),

$$\begin{aligned}
B_x(t_x) + \frac{t_x}{2} &= \int_0^{\frac{t_x}{2\alpha}} e^{-2xw^2} \sin[2wx + w(\alpha+2-2\beta)] \frac{dw}{w} \\
&= \int_0^{\frac{t_x}{2\alpha}} e^{-2xw^2} \sin[w(2x+\alpha+2-2\beta)] \frac{dw}{w}
\end{aligned}$$



The substitution  $z = (2x+2-2\beta+\alpha)w$ , yield

$$\int_0^{\left(\frac{x+1-\beta}{\alpha} + \frac{1}{2}\right)t_x} e^{-2xz^2/(2x+2-2\beta+\alpha)^2} \frac{\sin z}{z} dz$$

Hence for any positive number  $T$ , if  $x \rightarrow \infty$ ,  $t_x \rightarrow 0$  in such a way that  $xt_x \rightarrow \alpha T$ , then

$$B_x(t_x) = \int_0^T \frac{\sin z}{z} dz. \quad (2.5)$$

Thus equation (2.5) shows that  $(B, \alpha, \beta)$  transform preserves Gibbs phenomenon for Fourier series. When  $\alpha = \beta = 1$ , we obtain the result on Gibbs phenomenon for Borel summability proved by Lorch [30].

## Chapter III

### LEBESGUE CONSTANTS FOR $(B, \alpha, \beta)$ SUMMABILITY

In this chapter the Lebesgue constants for Borel-type method of summability is determined. These constants are defined as follows. Let  $D_n(t)$  the Dirichlet's kernel, namely  $\frac{\sin(2n+1)t}{\sin t}$  and  $B_x(t)$  the  $(B, \alpha, \beta)$  transform of the sequence  $\{D_n(t)\}$ . In finding the Lebesgue constants, we estimate the value of the following integral for large values of  $x$ ,

$$\frac{2}{\pi} \int_0^{\frac{\pi}{2}} |B_x(t)| dt \quad (3.1)$$

The following theorem generalizes the corresponding result for Borel summability [28].

**Theorem 3.1.** (c.f.[13])

If  $L_B(x)$  denote the Lebesgue constants for  $(B, \alpha, \beta)$  summability, then

$$L_B(x) = \frac{2}{\pi^2} \log\left(\frac{2x}{\pi^2}\right) - \frac{2}{\pi^2} C - \frac{2}{\pi^2} \int_0^{\frac{\pi}{2}} \psi\left(\frac{t}{\pi}\right) \sin t \, dt \\ + O\left(\frac{1}{\sqrt{x}}\right) \quad (x \rightarrow \infty)$$

where,

$$C = \int_0^1 \frac{1-e^{-y}}{y} dy - \int_1^{\infty} \frac{dy}{ye^y} \quad \text{is the Euler-Mascheroni constant and } \psi\left(\frac{t}{\pi}\right) = \frac{\Gamma'\left(\frac{t}{\pi}\right)}{\Gamma\left(\frac{t}{\pi}\right)}$$

Proof:

Let  $B_x(t)$  denote the  $(B, \alpha, \beta)$  transforms of the sequence of functions  $\left\{ \frac{\sin(2n+1)t}{\sin t} \right\}$

Then

$$\begin{aligned} B_x(t) &= \alpha e^{-x} \sum_{n=N}^{\infty} \frac{\sin(2n+1)t}{\sin t} \cdot \frac{x^{\alpha n + \beta - 1}}{\Gamma(\alpha n + \beta)} \\ &= \frac{\alpha e^{-x}}{\sin t} \operatorname{Im} \left[ \sum_{n=N}^{\infty} e^{i(2n+1)t} \cdot \frac{x^{\alpha n + \beta - 1}}{\Gamma(\alpha n + \beta)} \right] \\ &= \frac{\alpha e^{-x}}{\sin t} \operatorname{Im} \left[ e^{it - (\beta-1)\frac{i2t}{\alpha}} \sum_{n=N}^{\infty} \frac{(xe^{\frac{i2t}{\alpha}})^{\alpha n + \beta - 1}}{\Gamma(\alpha n + \beta)} \right] \\ &= \frac{\alpha e^{-x}}{\sin t} \operatorname{Im} \left[ e^{i\left\{t - \frac{2(\beta-1)t}{\alpha}\right\}} \cdot \frac{1}{\alpha} e^x e^{\frac{i2t}{\alpha}} \right] \\ &= \frac{e^{-x}}{\sin t} \operatorname{Im} \left[ e^{i\left\{t + \frac{2(1-\beta)}{\alpha}t + x(\sin \frac{2t}{\alpha})\right\}} \cdot e^x \cos \frac{2t}{\alpha} \right] \\ &= \frac{e^{-2x \sin^2 \frac{t}{\alpha}}}{\sin t} \cdot \sin \left[ x \sin \frac{2t}{\alpha} + 2t \left( \frac{1}{2} + \frac{1-\beta}{\alpha} \right) \right] \end{aligned}$$

It follows from (3.1) that the Lebesgue constants  $L_B(x)$  for  $(B, \alpha, \beta)$  method is given by

$$L_B(x) = \frac{2}{\pi} \int_0^{\frac{\pi}{2}} e^{-2x \sin^2 \frac{t}{\alpha}} \left| \sin \left[ x \sin \frac{2t}{\alpha} + 2t \left( \frac{1}{2} + \frac{1-\beta}{\alpha} \right) \right] \right| \frac{dt}{\sin t} \quad (3.2)$$

The evaluation of the integral is divided into many stages and hence the following lemmas.

Lemma 3.2.

Let,

$$L_{B_1}(x) = \frac{2}{\pi} \int_0^{\frac{\pi}{2}} e^{-2x \sin^2 \frac{t}{\alpha}} \left| \sin \left[ x \sin \frac{2t}{\alpha} + 2t \left( \frac{1}{2} + \frac{1-\beta}{\alpha} \right) \right] \right| \frac{dt}{t}$$

Then  $L_B(x) = L_{B_1}(x) + O\left(\frac{1}{\sqrt{x}}\right) \quad (x \rightarrow \infty)$

Proof:

Using the following inequalities

$$\begin{aligned} 0 < t - \sin t &< \frac{t^3}{3!} \quad \text{for } t \geq 0 \\ 1 < \frac{t}{\sin t} &< \frac{\pi}{2} \quad \text{for } 0 < t < \frac{\pi}{2} \end{aligned} \quad (3.3)$$

it follows that

$$\begin{aligned} 0 &\leq L_B(x) - L_{B_1}(x) \\ &= \frac{2}{\pi} \int_0^{\frac{\pi}{2}} \left[ \frac{1}{\sin t} - \frac{1}{t} \right] e^{-2x \sin^2 \frac{t}{\alpha}} \left| \sin \left[ x \sin \frac{2t}{\alpha} + 2t \left( \frac{1}{2} + \frac{1-\beta}{\alpha} \right) \right] \right| dt \end{aligned}$$

$$\begin{aligned}
&\ll \frac{2}{\pi} \int_0^{\frac{\pi}{2}} \frac{t - \sin t}{t \sin t} e^{-2x \sin^2 \frac{t}{\alpha}} dt \\
&\ll \frac{2}{\pi} \int_0^{\frac{\pi}{2}} \frac{t^3}{3!} \cdot \frac{1}{t} \cdot \frac{\pi}{2t} e^{-2x \sin^2 \frac{t}{\alpha}} dt \\
&= \frac{1}{6} \int_0^{\frac{\pi}{2}} t e^{-2x \sin^2 \frac{t}{\alpha}} dt
\end{aligned}$$

Substitution of  $y = \frac{t}{\alpha}$  reduces the integral to

$$\begin{aligned}
&\frac{\alpha^2}{6} \int_0^{\frac{\pi}{2\alpha}} y e^{-2x \sin^2 y} dy \\
&\ll \frac{\alpha^2}{6} \int_0^{M\pi} y e^{-2x \sin^2 y} dy \quad \text{where } M = \left[ \frac{1}{2\alpha} \right] + 1 \\
&= \frac{\alpha^2}{6} \sum_{n=1}^M \int_{(n-1)\pi}^{n\pi} y e^{-2x \sin^2 y} dy
\end{aligned}$$

Also,

$$\begin{aligned}
&\int_{(n-1)\pi}^{n\pi} y e^{-2x \sin^2 y} dy < n\pi \int_{(n-1)\pi}^{n\pi} e^{-2x \sin^2 y} dy \\
&= n\pi \int_0^{\pi} e^{-2x \sin^2 y} dy \\
&= 2n\pi \int_0^{\frac{\pi}{2}} e^{-2x \sin^2 y} dy \\
&\ll 2n\pi \int_0^{\frac{\pi}{2}} e^{-8x \frac{y^2}{\pi^2}} dy \quad (\text{using 3.3})
\end{aligned}$$

$$< 2n\pi \int_0^{\infty} e^{-8x \frac{y^2}{\pi^2}} dy$$

The substitution  $z = \sqrt{8x} \frac{y}{\pi}$  shows that

$$\int_0^{\infty} e^{-8x \frac{y^2}{\pi^2}} dy = O\left(\frac{1}{\sqrt{x}}\right) \quad (x \rightarrow \infty)$$

Since  $M$  is finite, it follows that

$$L_B(x) - L_{B_1}(x) = O\left(\frac{1}{\sqrt{x}}\right) \quad (x \rightarrow \infty)$$

Thus

$$L_B(x) = L_{B_1}(x) + O\left(\frac{1}{\sqrt{x}}\right) \quad (x \rightarrow \infty)$$

Lemma 3.3.

Let

$$L_{B_2}(x) = \frac{2}{\pi} \int_0^{\frac{\pi}{2}} e^{-2x \frac{t^2}{\alpha^2}} \left| \sin\left[x \sin \frac{2t}{\alpha} + 2t\left(\frac{1}{2} + \frac{1-\beta}{\alpha}\right)\right] \right| \frac{dt}{t}$$

Then

$$L_B(x) = L_{B_2}(x) + O\left(\frac{1}{\sqrt{x}}\right) \quad (x \rightarrow \infty)$$

Proof: We have

$$\begin{aligned} 0 &\leq L_{B_1}(x) - L_{B_2}(x) \\ &= \frac{2}{\pi} \int_0^{\frac{\pi}{2}} \left[ e^{-2x \sin^2 \frac{t}{\alpha}} - e^{-2x \frac{t^2}{\alpha^2}} \right] \\ &\quad \left| \sin\left[x \sin \frac{2t}{\alpha} + 2t\left(\frac{1}{2} - \frac{1-\beta}{\alpha}\right)\right] \right| \frac{dt}{t} \end{aligned}$$

$$= \frac{2}{\pi} \int_0^{\frac{\pi}{2\alpha}} [e^{-2x \sin^2 y} - e^{-2x y^2}] | \sin[x \sin 2y + 2\alpha y (\frac{1}{2} + \frac{1-\beta}{\alpha})] | \frac{dy}{y}$$

Consequently,

$$0 \leq L_{B_1}(x) - L_{B_2}(x) \leq \frac{2}{\pi} \int_0^{\frac{\pi}{2\alpha}} \frac{e^{2x(y^2 - \sin^2 y)} - 1}{e^{2xy^2} \cdot y} dy \quad (3.4)$$

For approximating the integral on the right hand side we consider the following cases for different values of  $\alpha$ .

#### Case-1

Let  $\alpha \geq \frac{1}{2}$ . This implies  $\frac{\pi}{2\alpha} \leq \pi$ .

From equation (3.4) it follows that,

$$\begin{aligned} L_{B_1}(x) - L_{B_2}(x) &\leq \frac{2}{\pi} \int_0^{\frac{\pi}{2\alpha}} \frac{e^{2x(y^2 - \sin^2 y)} - 1}{e^{2xy^2} \cdot y} dy \\ &= \frac{2}{\pi} \int_0^{\frac{\pi}{2}} \frac{e^{2x(y^2 - \sin^2 y)} - 1}{e^{2xy^2} \cdot y} dy + \\ &\quad \frac{2}{\pi} \int_{\frac{\pi}{2}}^{\frac{\pi}{2\alpha}} \frac{e^{2x(y^2 - \sin^2 y)} - 1}{e^{2xy^2} \cdot y} dy \\ &= I_1 + I_2, \end{aligned}$$

say. To estimate  $I_1$  and  $I_2$  we make use of the following inequalities,

$$\begin{aligned} \sin y &\leq y, \quad \text{for } y \geq 0 \\ e^y - 1 &< ye^y \text{ for } y \geq 0 \end{aligned} \quad (3.5)$$

Using the above inequality in  $I_1$ , we see that

$$\begin{aligned} I_1 &\leq \frac{2}{\pi} \int_0^{\frac{\pi}{2}} \frac{2x(y^2 - \sin^2 y) e^{2x(y^2 - \sin^2 y)}}{ye^{2xy^2}} \cdot dy \\ &= \frac{4x}{\pi} \int_0^{\frac{\pi}{2}} (y + \sin y)(y - \sin y) e^{-2x \sin^2 y} \cdot \frac{dy}{y} \\ &\leq \frac{4x}{\pi} \int_0^{\frac{\pi}{2}} 2y \cdot \frac{y^3}{3!} e^{-\frac{8x}{\pi^2} y^2} \cdot \frac{dy}{y} \quad (\text{Using (3.3)}) \\ &= \frac{4x}{3\pi} \int_0^{\frac{\pi}{2}} y^3 e^{-\frac{8x}{\pi^2} y^2} dy \end{aligned}$$

The substitution  $z = \frac{8x}{\pi^2} y^2$  shows that

$$I_1 \leq \frac{\pi^3}{96x} \int_0^{2x} ze^{-z} dz$$

Since the integral  $\int_0^{\infty} ze^{-z} dz$  converges, it is bounded.

It follows that

$$I_1 = O\left(\frac{1}{x}\right) \quad (x \rightarrow \infty)$$



Also,

$$\begin{aligned}
 I_2 &= \frac{2}{\pi} \int_{\frac{\pi}{2}}^{\pi} \frac{e^{2x(y^2 - \sin^2 y)} - 1}{e^{2xy^2} \cdot y} dy \\
 &\leq \frac{2}{\pi} \int_{\frac{\pi}{2}}^{\pi} e^{-2x \sin^2 y} \cdot \frac{dy}{y} \\
 &\leq \frac{4}{\pi^2} \int_{\frac{\pi}{2}}^{\pi} e^{-2x \sin^2 y} dy \\
 &= \frac{4}{\pi^2} \int_0^{\frac{\pi}{2}} e^{-2x \sin^2 y} dy \\
 &\leq \frac{4}{\pi^2} \int_0^{\frac{\pi}{2}} e^{\frac{-8x}{\pi^2} y^2} dy \quad (\text{using (3.3)}) \\
 &\leq \frac{4}{\pi^2} \int_0^{\infty} e^{\frac{-8x}{\pi^2} y^2} dy
 \end{aligned}$$

The substitution  $z = \sqrt{\frac{8x}{\pi}} y$  shows that

$$I_2 \leq \frac{4}{\pi\sqrt{8x}} \int_0^{\infty} e^{-z^2} dz = \frac{1}{\sqrt{2\pi x}}$$

Thus we see that

$$I_2 = O\left(\frac{1}{\sqrt{x}}\right) \quad (x \rightarrow \infty).$$

It follows that for  $\alpha \gg \frac{1}{2}$

$$L_{B_1}(x) - L_{B_2}(x) = O\left(\frac{1}{\sqrt{x}}\right) \quad (x \rightarrow \infty).$$

Case 2:

Let  $0 < \alpha < \frac{1}{2}$ . Set  $M = \left[ \frac{1}{2\alpha} \right] + 1$ .

Consequently from (3.4) we see that

$$\begin{aligned} L_{B_1}(x) - L_{B_2}(x) &\leq \frac{2}{\pi} \int_0^{M\pi} \frac{e^{2x(y^2 - \sin^2 y)} - 1}{ye^{2xy^2}} dy \\ &= \frac{2}{\pi} \int_0^{\pi} \frac{e^{2x(y^2 - \sin^2 y)} - 1}{ye^{2xy^2}} dy + \frac{2}{\pi} \int_{\pi}^{M\pi} \frac{e^{2x(y^2 - \sin^2 y)} - 1}{ye^{2xy^2}} dy \\ &= I_3 + I_4, \end{aligned}$$

say, As in case-1, we find that

$$I_3 = O\left(\frac{1}{\sqrt{x}}\right) \quad (x \rightarrow \infty).$$

Further,

$$I_4 = \frac{2}{\pi} \sum_{r=2}^M \int_{(r-1)\pi}^{r\pi} \frac{e^{2x(y^2 - \sin^2 y)} - 1}{ye^{2xy^2}} dy.$$

The following calculation shows that each integral on the right hand side is  $O\left(\frac{1}{\sqrt{x}}\right)$ .

To this end note that

$$\begin{aligned}
 \int_{(r-1)\pi}^{r\pi} \frac{e^{2x(y^2 - \sin^2 y)} - 1}{e^{2xy^2} \cdot y} dy &< \frac{1}{(r-1)\pi} \int_{(r-1)\pi}^{r\pi} e^{-2x \sin^2 y} dy \\
 &= \frac{2}{(r-1)\pi} \int_0^{\frac{\pi}{2}} e^{-2x \sin^2 y} dy \\
 &\leq \frac{2}{(r-1)\pi} \int_0^{\frac{\pi}{2}} e^{-\frac{8x}{\pi^2} y^2} dy \\
 &= O\left(\frac{1}{\sqrt{x}}\right) \quad (x \rightarrow \infty)
 \end{aligned}$$

Since  $M$  is finite, it follows that

$$I_4 = O\left(\frac{1}{\sqrt{x}}\right) \quad (x \rightarrow \infty)$$

Thus we see that for all values of  $\alpha$

$$L_{B_1}(x) - L_{B_2}(x) = O\left(\frac{1}{\sqrt{x}}\right) \quad (x \rightarrow \infty)$$

This result together with Lemma 3.2 shows that

$$L_B(x) = L_{B_2}(x) + O\left(\frac{1}{\sqrt{x}}\right) \quad (x \rightarrow \infty)$$

Lemma 3.4.

Let

$$L_{B_3}(x) = \frac{2}{\pi} \int_0^{\frac{\pi}{2}} e^{-2x \frac{t^2}{\alpha^2}} \left| \sin \left[ 2x \frac{t}{\alpha} + 2t \left( \frac{1}{2} + \frac{1-\beta}{\alpha} \right) \right] \right| \frac{dt}{t}$$

Then

$$L_B(x) = L_{B_3}(x) + O \left( \frac{1}{\sqrt{x}} \right) \quad (x \rightarrow \infty)$$

Proof: Now,

$$\begin{aligned} |L_{B_2}(x) - L_{B_3}(x)| &\leq \frac{2}{\pi} \int_0^{\frac{\pi}{2}} e^{-2x \frac{t^2}{\alpha^2}} \left| \sin \left[ x \sin \frac{2t}{\alpha} + \right. \right. \\ &\quad \left. \left. 2t \left( \frac{1}{2} + \frac{1-\beta}{\alpha} \right) \right] - \sin \left[ x \frac{2t}{\alpha} + 2t \left( \frac{1}{2} + \frac{1-\beta}{\alpha} \right) \right] \right| \frac{dt}{t} \end{aligned}$$

Applying the trigonometric formula

$$\sin C - \sin D = 2 \sin \frac{C-D}{2} \cos \frac{C+D}{2}$$

it follows that

$$\begin{aligned} |L_{B_2}(x) - L_{B_3}(x)| &\leq \frac{4}{\pi} \int_0^{\frac{\pi}{2}} e^{-2x \frac{t^2}{\alpha^2}} \left| \sin \left[ \frac{x}{2} \left( \frac{2t}{\alpha} - \sin \frac{2t}{\alpha} \right) \right] \right| \frac{dt}{t} \\ &\leq \frac{4x}{\pi} \int_0^{\frac{\pi}{2}} e^{-2x \frac{t^2}{\alpha^2}} \left( \frac{2t}{\alpha} - \sin \frac{2t}{\alpha} \right) \cdot \frac{dt}{t} \end{aligned}$$

$$\begin{aligned}
&\leq \frac{4x}{\pi} \int_0^{\frac{\pi}{2}} e^{-2x\frac{t^2}{\alpha^2}} \cdot \frac{8t^3}{6\alpha^3} \frac{dt}{2t} \quad (\text{using (3.3)}) \\
&= \frac{8x}{3\pi \alpha^3} \int_0^{\frac{\pi}{2}} t^2 e^{-2x\frac{t^2}{\alpha^2}} dt \\
&= \frac{8x}{3\pi} \int_0^{\frac{\pi}{2\alpha}} y^2 e^{-2xy^2} dy \\
&\leq \frac{8x}{3\pi} \int_0^{\infty} y^2 e^{-2xy^2} dy \\
&= \frac{8x}{3\pi} \left[ y \frac{e^{-2xy^2}}{-4x} \Big|_0^{\infty} + \int_0^{\infty} \frac{e^{-2xy^2}}{4x} dy \right] \\
&= \frac{2}{3\pi} \int_0^{\infty} e^{-2xy^2} dy = \frac{2}{3\pi\sqrt{2x}} \int_0^{\infty} e^{-z^2} dz
\end{aligned}$$

Thus  $L_{B_2}(x) - L_{B_3}(x) = O\left(\frac{1}{\sqrt{x}}\right)$  ( $x \rightarrow \infty$ )

Together with Lemma 3.3, it follows that

$$L_B(x) = L_{B_3}(x) + O\left(\frac{1}{\sqrt{x}}\right) \quad (x \rightarrow \infty)$$

Lemma 3.5.

$$\text{If } f_x(t) = \frac{1}{t} \left(1 - e^{-2x\frac{t^2}{\alpha^2}}\right)$$

then  $\int_0^{\frac{\pi}{2}} |f'_x(t)| dt = o(\sqrt{x}) \quad (x \rightarrow \infty)$

Proof:

Differentiating we obtain,

$$f'_x(t) = \frac{1}{t^2} \left[ \left( 4x \frac{t^2}{\alpha^2} + 1 \right) e^{-2x \frac{t^2}{\alpha^2}} - 1 \right]$$

set  $y = 2x \frac{t^2}{\alpha^2}$  and  $g(y) = t^2 f'_x(t)$  (3.6)

consequently,

$$g(y) = (2y+1) e^{-y} - 1$$

and

$$g(y) = 0 \iff 2y+1 = e^y$$

This clearly shows that, there exists a unique positive value for  $y$  say  $2\delta^2$  such that

$$\begin{aligned} g(2\delta^2) &= 0, \text{ and} \\ g(y) &< 0 \text{ if } y > 2\delta^2 \\ g(y) &> 0 \text{ if } y < 2\delta^2 \end{aligned} \quad (3.7)$$

(3.6) and (3.7) together shows that

$$f'_x(t) > 0 \text{ if } 0 < t < \frac{\alpha\delta}{\sqrt{x}}$$

and  $f'_x(t) < 0$  if  $t > \frac{\alpha\delta}{\sqrt{x}}$

consequently

$$\begin{aligned} \int_0^{\frac{\pi}{2}} |f_x'(t)| dt &= \int_0^{\frac{\alpha\delta}{\sqrt{x}}} f_x'(t) dt - \int_{\frac{\alpha\delta}{\sqrt{x}}}^{\frac{\pi}{2}} f_x'(t) dt. \\ &= f_x\left(\frac{\alpha\delta}{\sqrt{x}}\right) - f_x(0) - f_x\left(\frac{\pi}{2}\right) + f_x\left(\frac{\alpha\delta}{\sqrt{x}}\right) \\ &= \frac{2\sqrt{x}}{\alpha\delta} (1 - e^{-2\delta^2}) - \frac{2}{\pi} \left(1 - e^{-\frac{\pi^2 x}{2\alpha^2}}\right) \\ &= O(\sqrt{x}) \quad (x \rightarrow \infty) \end{aligned}$$

Lemma 3.6.

If

$$L(x) = \frac{2}{\pi} \int_0^{\frac{\pi}{2}} |\sin[2x \frac{t}{\alpha} + 2t(\frac{1}{2} + \frac{1-\beta}{\alpha})]| \frac{dt}{t}$$

then,

$$L(x) = \frac{4}{\pi^2} \log x - \frac{4}{\pi^2} \log \alpha - \frac{2}{\pi^2} \int_0^{\frac{\pi}{2}} \psi\left(\frac{t}{\pi}\right) \sin t \, dt + O\left(\frac{1}{x}\right) \quad (x \rightarrow \infty)$$

$$\text{where } \psi\left(\frac{t}{\pi}\right) = \left\lceil \left(\frac{t}{\pi}\right) \right\rceil / \left\lfloor \left(\frac{t}{\pi}\right) \right\rfloor$$

Proof:

$$L(x) = \frac{2}{\pi} \int_0^{\frac{\pi}{2}} |\sin [ \{ 2(\frac{x+1-\beta}{\alpha}) + 1 \} t ] | \frac{dt}{t} \quad (3.8)$$

To estimate this integral we make use of the following result proved in [29]

$$\text{If } L_1(x) = \frac{2}{\pi} \int_0^{\frac{\pi}{2}} \left| \frac{\sin(2x+1)t}{t} \right| dt$$

then,

$$L_1(x) = \frac{4}{\pi^2} \log x - \frac{2}{\pi^2} \int_0^{\pi} \psi\left(\frac{t}{\pi}\right) \sin t dt + O\left(\frac{1}{x}\right) \quad (x \rightarrow \infty)$$

It follows that

$$\begin{aligned} L(x) &= \frac{4}{\pi^2} \log \left( \frac{x+1-\beta}{\alpha} \right) - \frac{2}{\pi^2} \int_0^{\pi} \psi\left(\frac{t}{\pi}\right) \sin t dt + O\left(\frac{1}{\frac{x+1-\beta}{\alpha}}\right) \\ &\quad (x \rightarrow \infty) \\ &= \frac{4}{\pi^2} \log x - \frac{4}{\pi^2} \log \alpha - \frac{2}{\pi^2} \int_0^{\pi} \psi\left(\frac{t}{\pi}\right) \sin t dt + O\left(\frac{1}{x}\right) \\ &\quad (x \rightarrow \infty) \end{aligned}$$

Lemma 3.7. If

$$d(x) = L(x) - L_{B_3}(x)$$

then,

$$d(x) = \frac{2}{\pi^2} \log x + \frac{2}{\pi^2} \log \frac{\pi^2}{2\alpha^2} + \frac{2}{\pi^2} C + O\left(\frac{1}{\sqrt{x}}\right) \quad (x \rightarrow \infty)$$



where  $C$  is the Euler-Mascheroni constant given by

$$C = \int_0^1 \frac{1-e^{-y}}{y} dy - \int_1^{\infty} \frac{dy}{ye^y}$$

Proof:

Using the expressions of  $L(x)$  and  $L_{B_3}(x)$  from (3.8) and Lemma 3.4, we obtain,

$$\begin{aligned} d(x) &= \frac{2}{\pi} \int_0^{\frac{\pi}{2}} \left( \frac{1-e^{-2x\frac{t^2}{\alpha^2}}}{t} \right) \left| \sin \left[ \left\{ \frac{2(x+1-\beta)}{\alpha} + 1 \right\} t \right] \right| dt \\ &= \frac{4}{\pi^2} \int_0^{\frac{\pi}{2}} \left( \frac{1-e^{-2x\frac{t^2}{\alpha^2}}}{t} \right) dt + O\left(\frac{1}{\sqrt{x}}\right) \quad (x \rightarrow \infty) \end{aligned}$$

Set  $y = 2x\frac{t^2}{\alpha^2}$ . Then,

$$\begin{aligned} d(x) &= \frac{2}{\pi^2} \int_0^{\frac{\pi^2}{2\alpha^2}x} \frac{1-e^{-y}}{y} dy + O\left(\frac{1}{\sqrt{x}}\right) \quad (x \rightarrow \infty) \\ &= \frac{2}{\pi^2} \int_0^1 \frac{1-e^{-y}}{y} dy + \frac{2}{\pi^2} \int_1^{\frac{\pi^2}{2\alpha^2}x} \frac{1-e^{-y}}{y} dy + O\left(\frac{1}{\sqrt{x}}\right) \\ &= \frac{2}{\pi^2} \left[ \int_0^1 \frac{1-e^{-y}}{y} dy - \int_1^{\frac{\pi^2}{2\alpha^2}x} \frac{dy}{ye^y} \right] + \frac{2}{\pi^2} \log \left( \frac{\pi^2}{2\alpha^2}x \right) \\ &\quad + O\left(\frac{1}{\sqrt{x}}\right) \end{aligned}$$

$$\begin{aligned}
&= \frac{2}{\pi^2} \log x + \frac{2}{\pi^2} \log\left(\frac{\pi^2}{2\alpha^2}\right) + \frac{2}{\pi^2} \left[ \int_0^1 \frac{1-e^{-y}}{y} dy - \int_1^{\infty} \frac{dy}{ye^y} \right] \\
&\quad + \frac{2}{\pi^2} \int_{\frac{\pi^2}{2\alpha^2}x}^{\infty} \frac{dy}{ye^y} + O\left(\frac{1}{\sqrt{x}}\right)
\end{aligned}$$

The integral  $\int_{\frac{\pi^2}{2\alpha^2}x}^{\infty} \frac{dy}{ye^y}$  vanishes exponentially as  $x \rightarrow \infty$

and hence can be absorbed in  $O\left(\frac{1}{\sqrt{x}}\right)$ . Also it is known that the value of the integral

$$\int_0^1 \frac{1-e^{-y}}{y} dy - \int_1^{\infty} \frac{dy}{ye^y}$$

is equal to Euler-Mascherioni constant  $C$ .

Thus we see that

$$\begin{aligned}
d(x) &= L(x) - L_{B_3}(x) \\
&= \frac{2}{\pi^2} \log x + \frac{2}{\pi^2} \log\left(\frac{\pi^2}{2\alpha^2}\right) + \frac{2}{\pi^2} C + O\left(\frac{1}{\sqrt{x}}\right) \\
&\hspace{20em} (x \rightarrow \infty)
\end{aligned}$$

Proof of theorem.

Lemma 3.4, Lemma 3.6 and Lemma 3.7 together show that

$$\begin{aligned}L_B(x) &= L(x) - d(x) + O\left(\frac{1}{\sqrt{x}}\right) \quad (x \rightarrow \infty) \\&= \frac{2}{\pi^2} \log\left(\frac{2x}{\pi^2}\right) - \frac{2}{\pi^2} C - \frac{2}{\pi^2} \int_0^{\pi} \psi\left(\frac{t}{\pi}\right) \sin t \, dt \\&\quad + O\left(\frac{1}{\sqrt{x}}\right) \quad (x \rightarrow \infty)\end{aligned}$$

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## Chapter IV

### (B, $\alpha, \beta$ ) SUMMABILITY OF LEGENDRE SERIES

In this chapter the domain of summability of a series of Legendre polynomials by (B,  $\alpha, \beta$ ) method is obtained. The corresponding result for the Borel exponential method [22] follows as a particular case. The series to be considered is the expansion of  $(t-z)^{-1}$  in terms of Legendre polynomials. To this end consider the Legendre polynomials of the first and second kind of  $n^{\text{th}}$  degree denoted by  $P_n(z)$  and  $Q_n(w)$  respectively. The Laplace integral representation of  $P_n(z)$  and  $Q_n(w)$  are given by (c.f [40] ).

$$P_n = P_n(z) = \frac{1}{\pi} \int_0^{\pi} \xi^n d\phi \quad (4.1)$$

and

$$Q_n = Q_n(w) = \int_0^{\infty} \tau^{-n-1} d \quad (4.2)$$

where,

$$\xi = \xi(\phi) = z + \sqrt{(z^2-1)} \cos\phi$$

$$\tau = \tau(\psi) = w + \sqrt{(w^2-1)} \cosh\psi$$

The branch of  $\sqrt{(z^2-1)}$  is so chosen that  $z + \sqrt{(z^2-1)}$  lies

in the exterior of the unit circle.

Write,

$$s_k = s_k(z, w) = \sum_{n=0}^k (2n+1) P_n(z) Q_n(w)$$

and

$$d_n = d_n(z, w) = P_{n+1}(z) Q_n(w) - P_n(z) Q_{n+1}(w) \quad (4.3)$$

We have by Christoffel formula ([40], page 321)

$$\frac{1}{w-z} = s_n + (n+1) \frac{1}{w-z} d_n \quad (4.4)$$

By Hein's theorem ([40], page 322) the sequence  $\{s_n(z, w)\}$  converges to  $(w-z)^{-1}$  in the interior of the ellipse with foci  $\pm 1$  and passing through  $w$ . The following theorem asserts that the sequence  $\{s_n(z, w)\}$  is summable by  $(B, \alpha, \beta)$  method to  $(w-z)^{-1}$  in a wider region.

Theorem 4.1. (c.f [14] )

The sequence  $\{s_k(z, w)\}$  of partial sums of the series of Legendre polynomials  $\sum_{n=0}^{\infty} (2n+1) P_n(z) Q_n(w)$  is summable  $(B, \alpha, \beta)$  to  $(w-z)^{-1}$  in the region

$$\left\{ z: \operatorname{Re} \left\{ \frac{\xi(\phi)}{\tau(\psi)} \right\}^{\frac{1}{\alpha}} < \lambda, \left| \frac{\xi(\phi)}{\tau(\psi)} \right| < M, 0 \leq \phi \leq \pi, 0 \leq \psi \right\}$$

where  $M$  is a positive number and  $0 < \lambda < 1$ .

Proof:

Let  $T(x) = \sum_{k=N}^{\infty} u_k(x) s_k$ , where

$$u_k(x) = \alpha e^{-x} \frac{x^{\alpha k + \beta - 1}}{\Gamma(\alpha k + \beta)}$$

Then, using (3.4) we see that

$$\begin{aligned} T(x) &= \sum_{k=N}^{\infty} u_k(x) \left[ \frac{1}{w-z} - (k+1) \frac{1}{w-z} d_k \right] \\ &= \frac{1}{w-z} \sum_{k=N}^{\infty} u_k(x) - \frac{1}{w-z} \sum_{k=N}^{\infty} (k+1) u_k(x) d_k \\ &= \Sigma_1 - \Sigma_2 \end{aligned}$$

say. Applying Lemma 2.1, it follows that

$$\Sigma_1 = \frac{1}{w-z} [1 + o(1)] \quad (x \rightarrow \infty).$$

Hence

$$\lim_{x \rightarrow \infty} T(x) = (w-z)^{-1}$$

if and only if

$$\sum_{k=N}^{\infty} (k+1) u_k(x) d_k = o(1) \quad (x \rightarrow \infty) \quad (4.5)$$

We now investigate the region where (4.5) is satisfied. Using the equations (4.1), (4.2), (4.3) and (4.4) we have

$$\sum_{k=N}^{\infty} (k+1) u_k(x) d_k = \frac{1}{\pi} \sum_{k=N}^{\infty} (k+1) u_k(x) \int_0^{\infty} \int_0^{\pi} \left(\frac{\xi}{\tau}\right)^k \left(\frac{\xi}{\tau} - \frac{1}{\tau^2}\right) d\vartheta d\psi \quad (4.6)$$

Change of order of integration and summation in (4.6) is permissible if

$$\left| \frac{\xi}{\tau} \right| < M < \infty$$

with this assumption we obtain from (4.6)

$$\begin{aligned} \sum_{k=N}^{\infty} (k+1) u_k(x) d_k &= \frac{1}{\pi} \iint_{00}^{\infty\pi} \left(\frac{\xi}{\tau} - \frac{1}{\tau^2}\right) \left[ \sum_{k=N}^{\infty} (k+1) u_k(x) \left(\frac{\xi}{\tau}\right)^k \right] d\vartheta d\psi \\ &= o(1) \quad (x \rightarrow \infty) \end{aligned}$$

provided

$$\sum_{k=N}^{\infty} (k+1) u_k(x) \left(\frac{\xi}{\tau}\right)^k = o(1) \quad (x \rightarrow \infty)$$

Now

$$\sum_{k=N}^{\infty} (k+1) u_k(x) \left(\frac{\xi}{\tau}\right)^k = \alpha e^{-x} \sum_{k=N}^{\infty} (k+1) \frac{x^{\alpha k + \beta - 1} \left[\left(\frac{\xi}{\tau}\right)^{\alpha}\right]^{\frac{1}{\alpha}}}{\Gamma(\alpha k + \beta) \left(\frac{\xi}{\tau}\right)^{\frac{\beta-1}{\alpha}}}$$

$$= \alpha e^{-x} \left(\frac{\xi}{\Gamma}\right)^{\frac{1-\beta}{\alpha}} \sum_{k=N}^{\infty} (k+1) \frac{\left[x\left(\frac{\xi}{\Gamma}\right)^{\frac{1}{\alpha}}\right]^{\alpha k + \beta - 1}}{[(\alpha k + \beta)]} \quad (4.7)$$

Write  $I(x) = \sum_{k=N}^{\infty} \frac{\left[x\left(\frac{\xi}{\Gamma}\right)^{\frac{1}{\alpha}}\right]^{\alpha k + \beta - 1}}{[(\alpha k + \beta)]}$

Then  $I'(x) = \sum_{k=N}^{\infty} (\alpha k + \beta - 1) \frac{\left[x\left(\frac{\xi}{\Gamma}\right)^{\frac{1}{\alpha}}\right]^{\alpha k + \beta - 2}}{[(\alpha k + \beta)]} \left(\frac{\xi}{\Gamma}\right)^{\frac{1}{\alpha}}$

$$= \frac{\alpha}{x} \sum_{k=N}^{\infty} \frac{k \left[x\left(\frac{\xi}{\Gamma}\right)^{\frac{1}{\alpha}}\right]^{\alpha k + \beta - 1}}{[(\alpha k + \beta)]} + \frac{\beta - 1}{x} I(x)$$

consequently right side of equation (4.7) reduces to

$$\alpha e^{-x} \left(\frac{\xi}{\Gamma}\right)^{\frac{1-\beta}{\alpha}} \left[ \left\{ I'(x) - \frac{\beta - 1}{x} I(x) \right\} \frac{x}{\alpha} + I(x) \right] \quad (4.8)$$

For large values of  $x$ , Lemma 2.1 shows that

$$I(x) = \frac{1}{\alpha} e^{x\left(\frac{\xi}{\Gamma}\right)^{\frac{1}{\alpha}}}$$

$$I'(x) = \frac{1}{\alpha} e^{x\left(\frac{\xi}{\Gamma}\right)^{\frac{1}{\alpha}}} \left(\frac{\xi}{\Gamma}\right)^{\frac{1}{\alpha}}$$

Consequently from (4.7) and (4.8) we obtain



$$\begin{aligned}
\sum_{k=N}^{\infty} (k+1) u_k(x) d_k &= \alpha e^{-x} \left(\frac{s}{\tau}\right)^{\frac{1-\beta}{\alpha}} \left[ \frac{x}{\alpha} I'(x) + \frac{\alpha+1-\beta}{\alpha} I(x) \right] \\
&= \alpha e^{-x} \left(\frac{s}{\tau}\right)^{\frac{1-\beta}{\alpha}} \left[ \frac{x}{\alpha} \cdot \frac{1}{\alpha} e^{x\left(\frac{s}{\tau}\right)^{\frac{1}{\alpha}}} \left(\frac{s}{\tau}\right)^{\frac{1}{\alpha}} + \frac{\alpha+1-\beta}{\alpha} \cdot \frac{1}{\alpha} e^{x\left(\frac{s}{\tau}\right)^{\frac{1}{\alpha}}} \right] \\
&= \frac{1}{\alpha} e^{-x} e^{x\left(\frac{s}{\tau}\right)^{\frac{1}{\alpha}}} \left[ \left(\frac{s}{\tau}\right)^{\frac{1-\beta}{\alpha}} [(\alpha+1-\beta) + x\left(\frac{s}{\tau}\right)^{\frac{1}{\alpha}}] \right] \quad (4.9)
\end{aligned}$$

Let  $\left(\frac{s}{\tau}\right)^{\frac{1}{\alpha}} = a+ib$ . Then right hand side of (4.9) is equal to

$$\begin{aligned}
&= e^{-x(1-a)} e^{ixb} \left[ \frac{\alpha+1-\beta+x(a+ib)}{\frac{\beta-1}{\alpha} \alpha(a+ib)} \right] \\
&= o(1) \quad (x \rightarrow \infty)
\end{aligned}$$

if  $1-a > 0$ . That is real part of

$$\left(\frac{s}{\tau}\right)^{\frac{1}{\alpha}} < \lambda < 1.$$

## Chapter-V

### TAUBERIAN CONSTANTS FOR BOREL-TYPE SUMMABILITY

In this chapter we obtain Tauberian constants for Borel-type summability. The corresponding result for Borel-summability proved in [1] follows as a particular case of the following theorem

Theorem 5.1. (c.f. [15])

Let  $\Sigma u_k$  satisfy the tauberian condition

$$\limsup | \sqrt{k} u_k | = L < \infty \quad (5.1)$$

Let  $m \rightarrow \infty$ ,  $t \rightarrow \infty$  such that

$$\limsup \frac{m-t}{\sqrt{t}} = Q < \infty \quad (5.2)$$

Then

$$\limsup |B(t) - s_m| < AL$$

where

$$A = \sqrt{\frac{2\alpha}{\pi}} Q \int_0^Q e^{-\frac{\alpha}{2}z^2} dz + \sqrt{\frac{2}{\alpha\pi}} e^{-\frac{\alpha}{2}Q^2}$$

Lemma 5.2. (c.f.[1]).

It is easily seen that

(i) If  $0 < k < \mu$ , then

$$\sum_{j=k+1}^{\mu} \frac{1}{j^2} = \int_k^{\mu} \frac{1}{x^2} dx - \epsilon_k = 2(\sqrt{\mu} - \sqrt{k}) - \epsilon_k$$

where  $\epsilon_k \rightarrow 0$  as  $k \rightarrow \infty$ .

$$(ii) \quad |\sqrt{m} - \sqrt{\mu}| \frac{1}{m^{\frac{1}{2}}} < |m - \mu| \quad (m, \mu > 0)$$

Proof of the theorem:

Let  $j^* = \max(j, N)$ ,  $\{s_k\}$  the sequence of partial sums of  $\sum_0^{\infty} u_j$  and  $B(t)$  the  $(B, \alpha, \beta)$  transform of  $\{s_k\}$ . Then

$$\begin{aligned} B(t) &= \alpha e^{-t} \sum_{k=N}^{\infty} \frac{t^{\alpha k + \beta - 1}}{\Gamma(\alpha k + \beta)} \left( \sum_{j=0}^k u_j \right) \\ &= \sum_{j=0}^{\infty} \left[ \alpha e^{-t} \sum_{k=j^*}^{\infty} \frac{t^{\alpha k + \beta - 1}}{\Gamma(\alpha k + \beta)} \right] u_j \end{aligned}$$

Consequently

$$\begin{aligned}
 B(t) - s_m &= \sum_{j=0}^m \left[ - \left\{ 1 - \alpha e^{-t} \sum_{j^*}^{\infty} \frac{t^{\alpha k + \beta - 1}}{\Gamma(\alpha k + \beta)} \right\} \right] u_j \\
 &\quad + \sum_{j=m+1}^{\infty} \left[ \alpha e^{-t} \sum_{j^*}^{\infty} \frac{t^{\alpha k + \beta - 1}}{\Gamma(\alpha k + \beta)} \right] u_j \\
 &= \sum_{j=0}^m \left[ - \alpha e^{-t} \sum_N^{j^*-1} \frac{t^{\alpha k + \beta - 1}}{\Gamma(\alpha k + \beta)} \right] u_j \\
 &\quad + \sum_{j=m+1}^{\infty} \left[ \alpha e^{-t} \sum_{j^*}^{\infty} \frac{t^{\alpha k + \beta - 1}}{\Gamma(\alpha k + \beta)} \right] u_j
 \end{aligned}$$

$$\begin{aligned}
 \text{Let } c_{\mu} &= - \alpha \mu^{\frac{-1}{2}} e^{-t} \sum_N^{\mu-1} \frac{t^{\alpha k + \beta - 1}}{\Gamma(\alpha k + \beta)} \text{ if } N \leq \mu \leq m \\
 &= \alpha \mu^{\frac{-1}{2}} \sum_{\mu}^{\infty} \frac{t^{\alpha k + \beta - 1}}{\Gamma(\alpha k + \beta)} \text{ if } \mu > m
 \end{aligned}$$

Let  $x_k = \sqrt{k} u_k$  so that

$$\limsup |x_k| = L < \infty$$

and

$$B(t) - s_m = \sum c_\mu x_\mu$$

For fixed  $\mu$ ,  $\lim_{t \rightarrow \infty} c_\mu = 0$ .

To determine the Tauberian constant  $A$  in the theorem we follow the method used in [1] where it is shown that

$$A = \limsup \sum |c_\mu| \quad (5.3)$$

Use of Lemma 5.2 shows that

$$\begin{aligned} \sum_N^\infty |c_\mu| &= \alpha e^{-t} \sum_{\mu=N}^{m-1} \frac{t^{\alpha\mu+\beta-1}}{\Gamma(\alpha\mu+\beta)} \left( \sum_{\mu+1}^m j^{\frac{-1}{2}} \right) \\ &\quad + \alpha e^{-t} \sum_{\mu=m+1}^\infty \frac{t^{\alpha\mu+\beta-1}}{\Gamma(\alpha\mu+\beta)} \left( \sum_{m+1}^{\mu} j^{\frac{-1}{2}} \right) \\ &= 2\alpha e^{-t} \sum_{\mu=N}^{m-1} \frac{t^{\alpha\mu+\beta-1}}{\Gamma(\alpha\mu+\beta)} (\sqrt{m} - \sqrt{\mu}) \\ &\quad + 2\alpha e^{-t} \sum_{\mu=m+1}^\infty \frac{t^{\alpha\mu+\beta-1}}{\Gamma(\alpha\mu+\beta)} (\sqrt{\mu} - \sqrt{m}) \\ &\quad - \alpha e^{-t} \sum_{\mu=N}^{m-1} \frac{t^{\alpha\mu+\beta-1}}{\Gamma(\alpha\mu+\beta)} \epsilon_\mu - \epsilon_m \alpha e^{-t} \sum_{\mu=m+1}^\infty \frac{t^{\alpha\mu+\beta-1}}{\Gamma(\alpha\mu+\beta)} \end{aligned}$$

Since  $\epsilon_\mu$  and  $\epsilon_m$  are null sequences and  $(B, \alpha, \beta)$  transform is regular it follows that

$$\sum_N^\infty |c_\mu| = 2\alpha e^{-t} \sum_N^\infty \frac{t^{\alpha\mu+\beta-1}}{\Gamma(\alpha\mu+\beta)} |V_m - V_\mu| + o(1) \quad (m \rightarrow \infty)$$

For each  $\delta > 0$ , define

$$B_{\frac{1}{2}} = \left\{ \mu / |\mu - t| > \delta \sqrt{t} \right\}$$

Then

$$\begin{aligned} \sum_{\mu \in N}^\infty |c_\mu| &= 2\alpha e^{-t} \sum_{\mu \in B_{\frac{1}{2}}} \frac{t^{\alpha\mu+\beta-1}}{\Gamma(\alpha\mu+\beta)} |V_m - V_\mu| \\ &\quad + 2\alpha e^{-t} \sum_{\mu \notin B_{\frac{1}{2}}} \frac{t^{\alpha\mu+\beta-1}}{\Gamma(\alpha\mu+\beta)} |V_m - V_\mu| \\ &= \Sigma_1 + \Sigma_2, \text{ say.} \end{aligned}$$

Use of Lemma 5.2 (ii) and Cauchy-Schwartz inequality on  $\Sigma_1$  shows that

$$\Sigma_1 \leq \frac{2}{\sqrt{m}} \alpha e^{-t} \sum_{\mu \in B_{\frac{1}{2}}} \frac{t^{\alpha\mu+\beta-1}}{\Gamma(\alpha\mu+\beta)} |m - \mu|$$

$$\ll \frac{2}{\sqrt{m}} \left[ \alpha e^{-t} \sum_{\mu \in B_{\frac{1}{2}}} (m-\mu)^2 \frac{t^{\alpha\mu+\beta-1}}{\Gamma(\alpha\mu+\beta)} \right]^{1/2} \left[ \alpha e^{-t} \sum_{\mu \in B_{\frac{1}{2}}} \frac{t^{\alpha\mu+\beta-1}}{\Gamma(\alpha\mu+\beta)} \right]^{1/2}$$

An application of Lemma 2.2 shows that  $\Sigma_1 \rightarrow 0$ .

Now we shall consider the behaviour of the second sum

$$\Sigma_2 = 2\alpha e^{-t} \sum_{\mu \notin B_{\frac{1}{2}}} \frac{t^{\alpha\mu+\beta-1}}{\Gamma(\alpha\mu+\beta)} |\sqrt{m} - \sqrt{\mu}|$$

$$= 2\alpha e^{-t} \sum_{|\mu-t| \leq \delta\sqrt{t}} \frac{t^{\alpha\mu+\beta-1}}{\Gamma(\alpha\mu+\beta)} |\sqrt{m} - \sqrt{\mu}|$$

Condition (5.2) in the theorem implies that

$$\frac{m-t}{t} \rightarrow 0 \text{ and hence } \frac{m}{t} \rightarrow 1.$$

Also  $|\mu-t| \leq \delta\sqrt{t}$  implies

$$\left| \frac{\mu}{m} - \frac{t}{m} \right| < \frac{\delta\sqrt{t}}{m}$$

It follows that if  $\mu \notin B_{\frac{1}{2}}$ ,  $\frac{\mu}{m} \rightarrow 1$ .

Consequently

$$|\sqrt{m} - \sqrt{\mu}| = \frac{|m-\mu|}{\sqrt{m}+\sqrt{\mu}} \simeq \frac{|m-\mu|}{2\sqrt{m}}$$

Thus we see that

$$\begin{aligned} \Sigma_2 &= \frac{\alpha e^{-t}}{\sqrt{m}} \mu \notin B_{\frac{1}{2}} \sum \frac{t^{\alpha\mu+\beta-1}}{\Gamma(\alpha\mu+\beta)} |m-\mu| \\ &= \frac{\alpha e^{-t}}{\sqrt{m}} \sum_{\mu=N}^{\infty} \frac{t^{\alpha\mu+\beta-1}}{\Gamma(\alpha\mu+\beta)} |m-\mu| + o(1) \end{aligned}$$

Let

$$\begin{aligned} \phi(t) &= \frac{\alpha e^{-t}}{\sqrt{m}} \sum_{\mu=N}^{\infty} \frac{t^{\alpha\mu+\beta-1}}{\Gamma(\alpha\mu+\beta)} |m-\mu| \\ &= \frac{\alpha e^{-t}}{\sqrt{m}} \sum_{\mu=N}^m \frac{t^{\alpha\mu+\beta-1}}{\Gamma(\alpha\mu+\beta)} (m-\mu) + \frac{\alpha e^{-t}}{\sqrt{m}} \sum_{\mu=m+1}^{\infty} \frac{t^{\alpha\mu+\beta-1}}{\Gamma(\alpha\mu+\beta)} (\mu-m) \\ &= \Sigma_3 + \Sigma_4, \text{ say.} \end{aligned}$$

Now,

$$\begin{aligned} \Sigma_3 &= \frac{\alpha e^{-t}}{\sqrt{m}} \sum_{\mu=N}^m \frac{t^{\alpha\mu+\beta-1}}{\Gamma(\alpha\mu+\beta)} [(m-t) - (\mu-t)] \\ &= \Sigma_{31} - \Sigma_{32}, \text{ say} \end{aligned}$$



Then,

$$\Sigma_{31} = \frac{m-t}{\sqrt{m}} \alpha e^{-t} \sum_{\mu=N}^m \frac{t^{\alpha\mu+\beta-1}}{(\alpha\mu+\beta)}$$

Let  $\frac{1}{2} < \xi < \frac{2}{3}$ . Use of Lemma 2.2 shows that

$$\begin{aligned} \Sigma_{31} &= \frac{m-t}{\sqrt{m}} \sum_{t-t\xi}^m \sqrt{\frac{\alpha}{2\pi t}} e^{\frac{-\alpha}{2t}(\mu-t)^2} + o(1) \\ &= \frac{m-t}{\sqrt{m}} \sqrt{\frac{\alpha}{2\pi t}} \int_{t-t\xi}^m e^{\frac{-\alpha}{2t}(x-t)^2} dx + o(1) \end{aligned}$$

The substitution  $z = \frac{x-t}{\sqrt{t}}$  and the assumption (5.2) in the theorem show that

$$\begin{aligned} \Sigma_{31} &= \sqrt{\frac{\alpha}{2\pi}} Q \int_{-\infty}^Q e^{\frac{-\alpha}{2} z^2} dz + o(1) \\ &= o(1) + \sqrt{\frac{\alpha}{2\pi}} Q \left[ \int_{-\infty}^0 e^{\frac{-\alpha}{2} z^2} dz + \int_0^Q e^{-\frac{\alpha}{2} z^2} dz \right] \\ &= o(1) + \sqrt{\frac{\alpha}{2\pi}} Q \left[ \sqrt{\frac{\pi}{2\alpha}} + \int_0^Q e^{-\frac{\alpha}{2} z^2} dz \right] \\ &= o(1) + \frac{Q}{2} + \sqrt{\frac{\alpha}{2\pi}} Q \int_0^Q e^{-\frac{\alpha}{2} z^2} dz. \end{aligned}$$

A similar treatment of  $\Sigma_{32}$  yield

$$\begin{aligned}\Sigma_{32} &= \frac{\alpha e^{-t}}{\sqrt{m}} \sum_{\mu=N}^{\infty} \frac{t^{\alpha\mu+\beta-1}}{\Gamma(\alpha\mu+\beta)} (\mu-t) \\ &= \frac{1}{\sqrt{m}} \sqrt{\frac{\alpha}{2\pi t}} \int_{t-t^{\frac{1}{m}}}^m (x-t) e^{-\frac{\alpha}{2}(x-t)^2} dx + o(1)\end{aligned}$$

The substitution  $z = \frac{x-t}{\sqrt{t}}$  together with the fact  $\lim \frac{m}{t} = 1$  shows that

$$\begin{aligned}\Sigma_{32} &= \sqrt{\frac{\alpha}{2\pi}} \int_{-\infty}^Q z e^{-\frac{\alpha}{2} z^2} dz + o(1) \\ &= o(1) + \sqrt{\frac{\alpha}{2\pi}} \int_{-\infty}^{-Q} z e^{-\frac{\alpha}{2} z^2} dz \\ &= o(1) - \sqrt{\frac{1}{2\alpha\pi}} e^{-\frac{\alpha}{2} Q^2}\end{aligned}$$

Thus we see that

$$\begin{aligned}\Sigma_3 &= \Sigma_{31} - \Sigma_{32} \\ &= \frac{Q}{2} + \sqrt{\frac{\alpha}{2\pi}} Q \int_0^Q e^{-\frac{\alpha}{2} z^2} dz + \sqrt{\frac{1}{2\alpha\pi}} e^{-\frac{\alpha}{2} Q^2}\end{aligned}$$

A similar calculation with  $\Sigma_4$  will show that

$$\Sigma_4 = \Sigma_{41} - \Sigma_{42}, \text{ where}$$

$$\Sigma_{41} = \sqrt{\frac{\alpha}{2\pi}} \int_Q^\infty e^{-\frac{\alpha}{2}z^2} \cdot z \, dz = \frac{1}{\sqrt{2\alpha\pi}} e^{-\frac{\alpha}{2}Q^2}$$

$$\Sigma_{42} = \sqrt{\frac{\alpha}{2\pi}} Q \int_Q^\infty e^{-\frac{\alpha}{2}z^2} \, dz = \frac{Q}{2} - \sqrt{\frac{\alpha}{2\pi}} Q \int_0^Q e^{-\frac{\alpha}{2}z^2} \, dz$$

Thus we see that

$$\begin{aligned} \sum_{\mu=N}^{\infty} |c_\mu| &< \Sigma_3 + \Sigma_4 \\ &= \sqrt{\frac{2\alpha}{\pi}} Q \int_0^Q e^{-\frac{\alpha}{2}z^2} \, dz + \sqrt{\frac{2}{\alpha\pi}} e^{-\frac{\alpha}{2}Q^2} \end{aligned} \quad (5.4)$$

(5.3) and (5.4) together completes the proof of the theorem.

When  $\alpha=\beta=1$ , the value reduces to

$$\sqrt{\frac{2}{\pi}} \left[ Q \int_0^Q e^{-\frac{1}{2}z^2} \, dz + e^{-\frac{Q^2}{2}} \right]$$

which is proved in [1].

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